

Quantum info of LQG states



Pietro Dona



Collaboration with Eugenio Bianchi
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and w.i.p.

Plan of the talk:

Motivations and background

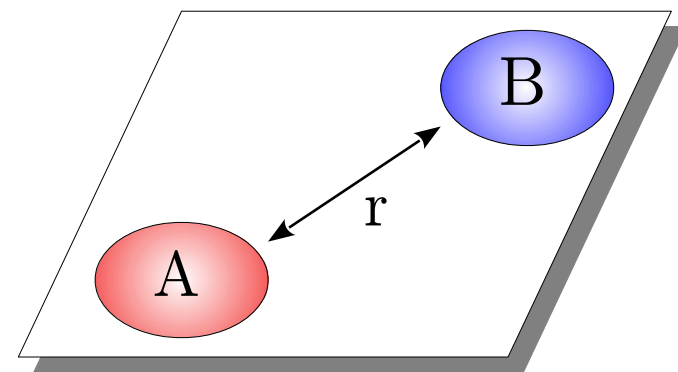
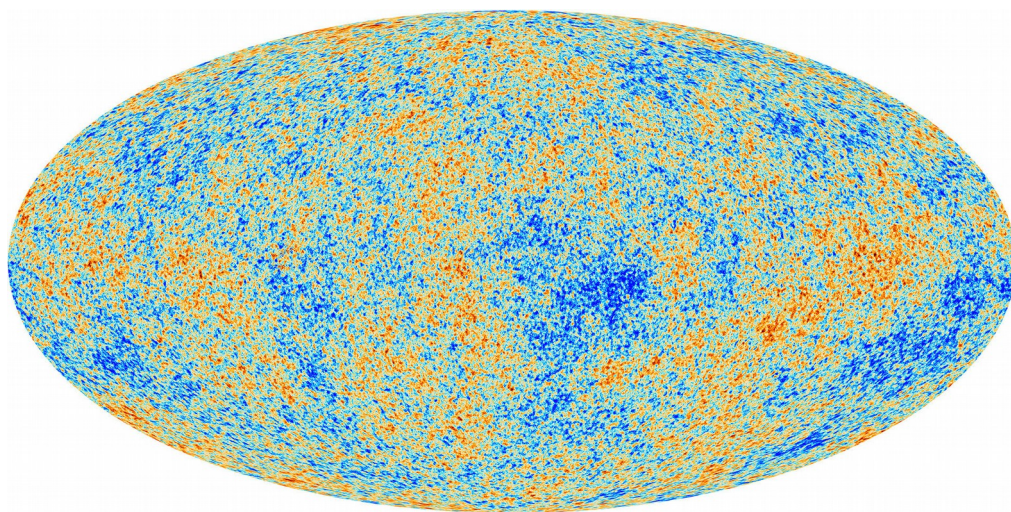
Typical entropy: Page curve and its variance

Entropy and correlations in Bell-network states

Correlations at space-like separation

Quantum field theory \longrightarrow Fock space contains:

- ~~(i) states with no space-like correlations~~
- (ii) states *with* specific short-ranged correlations



(e.g. Minkowski vacuum)

$$\langle \phi(0)\phi(r) \rangle \approx r^{-2}$$

Loop quantum gravity

- (i) states with no space-like correlations (spin-networks)
- (ii) states *with* specific short-ranged correlations (many proposals)

$$\mathcal{H} = \bigoplus_{\Gamma} \bigoplus_{j_{\ell}} \bigotimes_n \mathcal{H}_n$$

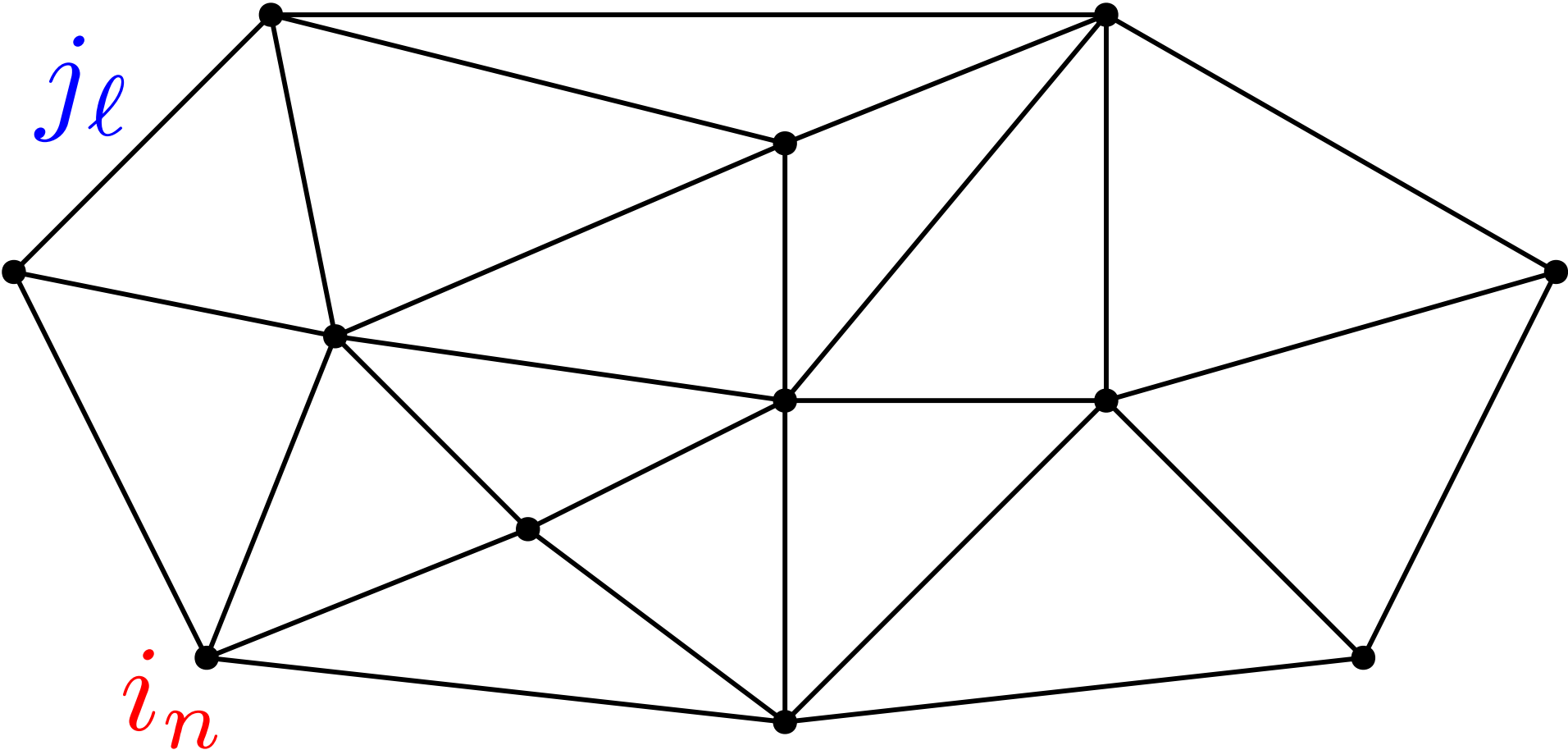
LQG Hilbert space

$$\mathcal{H}_\Gamma = \bigoplus_{j_\ell} \bigotimes_n \mathcal{H}_n$$



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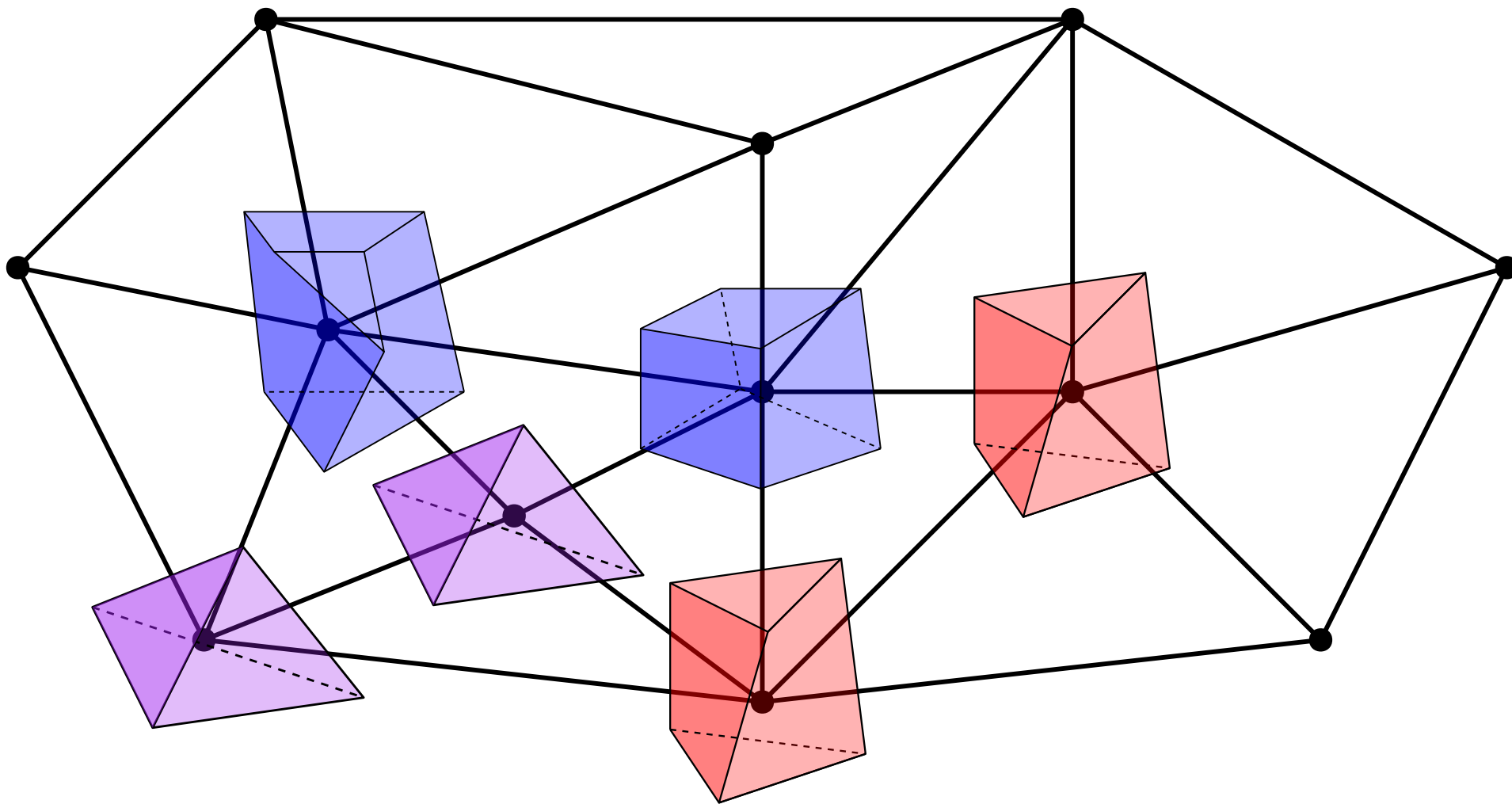


$$\mathcal{H}_n = \text{Inv} (j_1 \otimes \cdots \otimes j_F)$$

Quantum Polyhedra

[Bianchi, P.D., Speziale PRD 2010]

$$\mathcal{H}_\Gamma = \bigoplus_{j_\ell} \bigotimes_n \mathcal{H}_n$$



Geometric picture from LQG states

Information-theoretic bounds on correlations

State

$$|\psi\rangle \in \mathcal{H}_\Gamma$$

Subsystem A

$$O_A \in \mathcal{A}_A$$

Correlations

$$\mathcal{C}(O_A, O_B) = \langle \psi | O_A O_B | \psi \rangle - \langle \psi | O_A | \psi \rangle \langle \psi | O_B | \psi \rangle$$

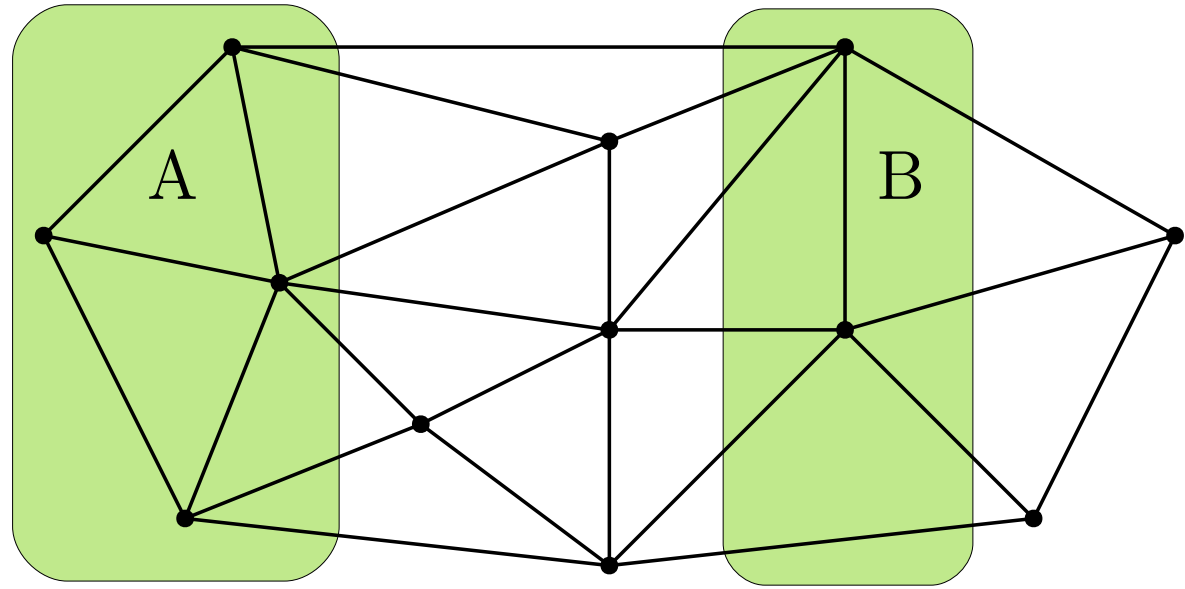
Bounded by mutual information

$$\frac{1}{2} \left(\frac{\mathcal{C}(O_A, O_B)}{\|O_A\| \|O_B\|} \right)^2 \leq S_A(\psi) + S_B(\psi) - S_{AB}(\psi)$$

[Wolf, Verstraete, Hastings, Cirac PRL 2008]

Entropy zero law or volume law = no correlations

What about a random state?



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Typical entropy: Page curve and its variance

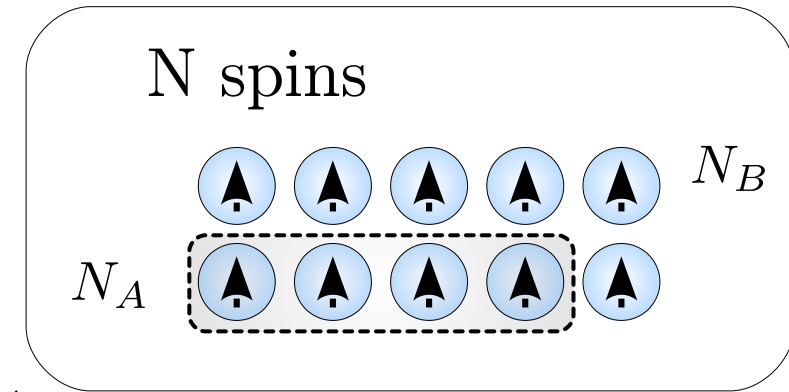
Entropy and correlations in Bell-network states

Typical entropy of a subsystem (with an example)

Hilbert space $\mathcal{H} = \bigotimes_{i=1}^N \mathbb{C}^2$ $d = 2^N$

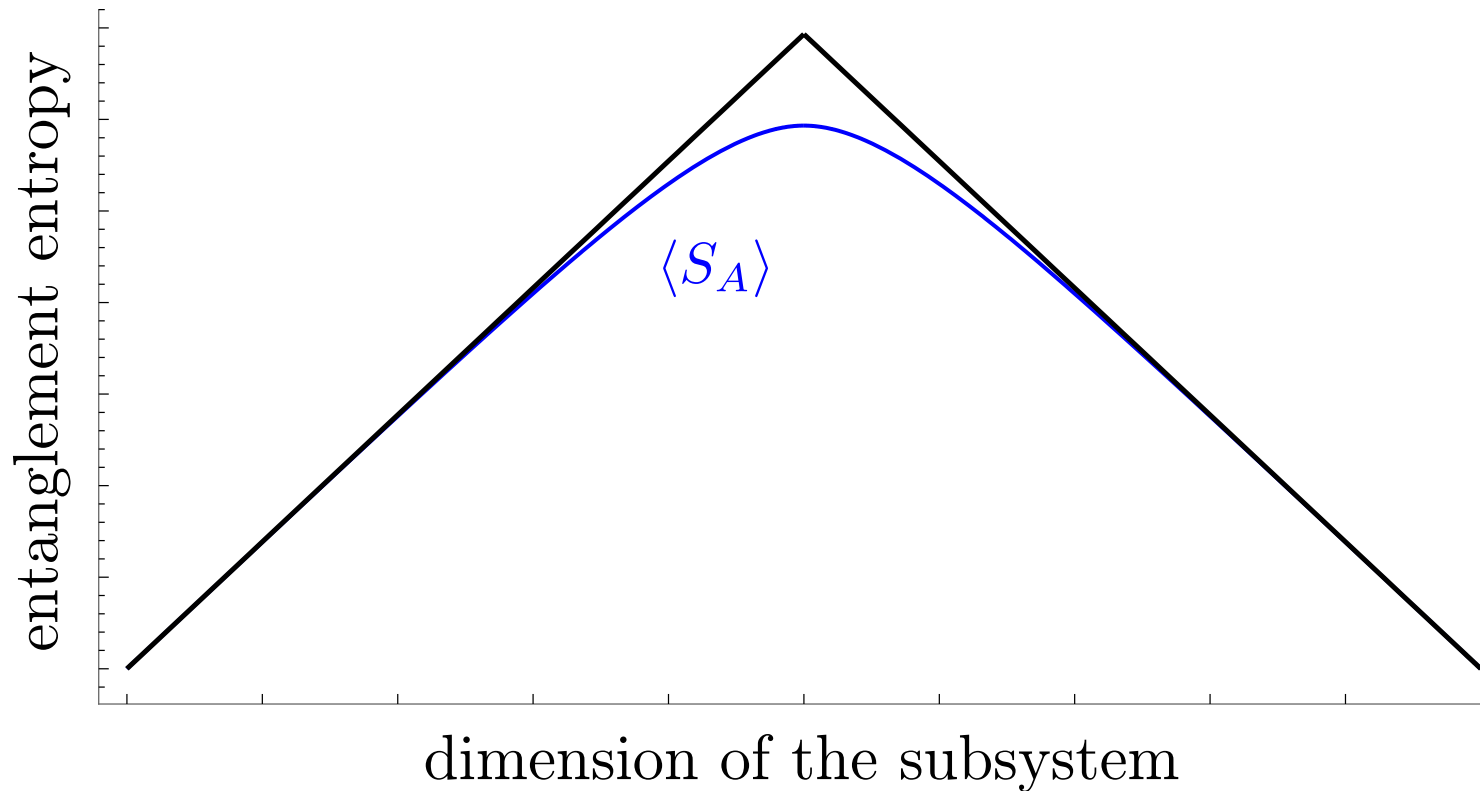
Subsystem N_A spins $d_A = 2^{N_A}$

$\rho_A = \text{Tr}_B (|\psi\rangle \langle\psi|)$ $S_A(\psi) = -\text{Tr}_A \rho_A \log \rho_A$



The average entropy (random uniform) is close to maximal

[Page PRL 1995]

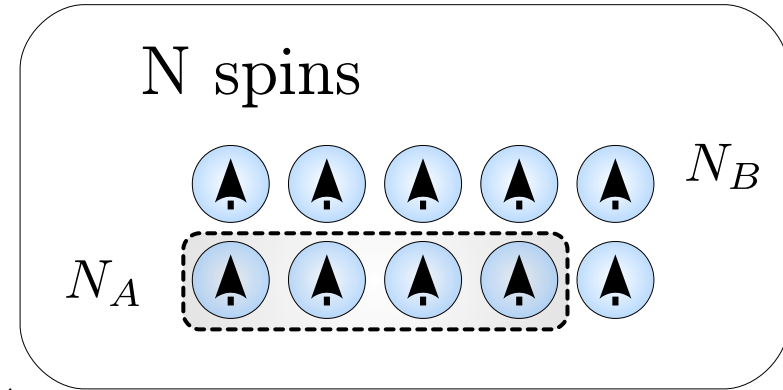


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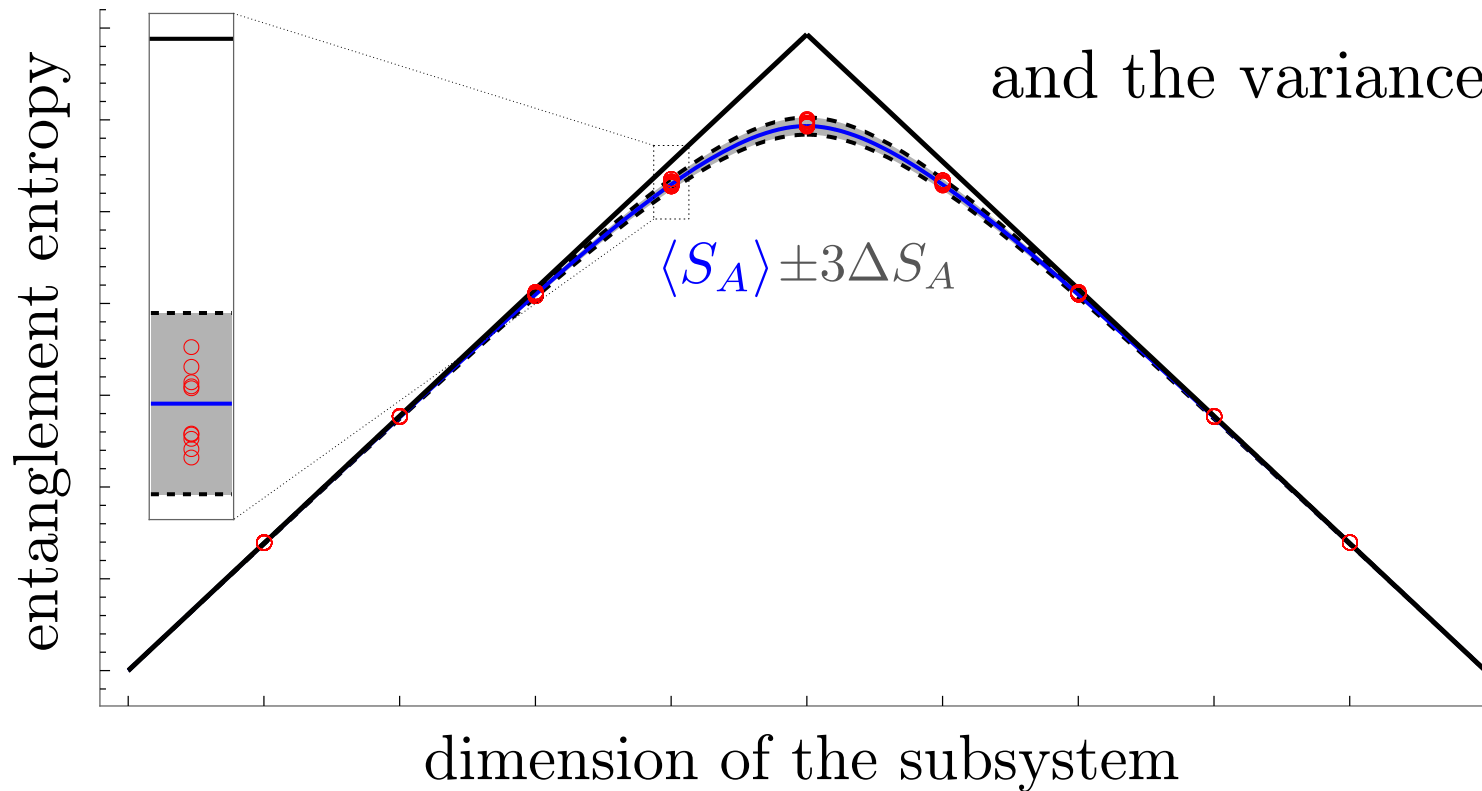


The average entropy (random uniform) is close to maximal

[Page PRL 1995]

and the variance is small

[Bianchi, P.D. PRD 2019]



Sketch of the proof

step 1:

$$\langle \text{Tr} \rho_A^r \rangle = \int d\mu(\psi) \text{Tr} \rho_A^r = \int \left(\sum_{a=1}^{d_A} \lambda_a^r \right) \mu(\lambda_1, \dots, \lambda_{d_A}) \prod_{b=1}^{d_A} d\lambda_b$$

step 2:

$$\langle S_A \rangle = - \langle \text{Tr} \rho_A \log \rho_A \rangle = - \lim_{r \rightarrow 1} \partial_r \langle \text{Tr} \rho_A^r \rangle$$

Eigenvalues of the density matrix
and induced integration measure
[Lloyd, Pagels Ann. Phys 1988]

exact result:

$$\langle S_A \rangle = \Psi(d_A d_B + 1) - \Psi(d_B + 1) - \frac{d_A - 1}{2d_B}$$

digamma function $\Psi = \Gamma'/\Gamma$

asymptotic:

$$\langle S_A \rangle \approx \log d_A - \frac{d_A^2 - 1}{2d_A d_B} \quad d_B \gg 1$$

$$d_A \leq d_B$$

Sketch of the proof

[Bianchi, P.D. PRD 2019]

variance:

$$\Delta S_A^2 = \langle S_A^2 \rangle - \langle S_A \rangle^2$$

step 1:

$$\langle \text{Tr} \rho_A^{r_1} \text{Tr} \rho_A^{r_2} \rangle = \int \left(\sum_{a=1}^{d_A} \lambda_a^{r_1} \right) \left(\sum_{a=1}^{d_A} \lambda_a^{r_2} \right) \mu(\lambda_1, \dots, \lambda_{d_A}) \prod_{b=1}^{d_A} d\lambda_b$$

step 2:

$$\langle S_A^2 \rangle = \lim_{\substack{r_1 \rightarrow 1 \\ r_2 \rightarrow 1}} \partial_{r_1} \partial_{r_2} \langle \text{Tr} \rho_A^{r_1} \text{Tr} \rho_A^{r_2} \rangle$$

exact result:

$$\Delta S_A^2 = \frac{d_A + d_B}{d_A d_B + 1} \Psi'(d_B + 1) - \Psi'(d_A d_B + 1) - \frac{(d_A - 1)(d_A + 2d_B - 1)}{4d_B^2(d_A d_B + 1)}$$

asymptotic:

$$\Delta S_A \approx \frac{1}{d_A d_B} \sqrt{\frac{d_A^2 - 1}{2}} \quad d_B \gg 1 \quad \Delta S_A \ll 1$$

The average entropy is typical!

$$d_A \leq d_B$$

10^5 samples

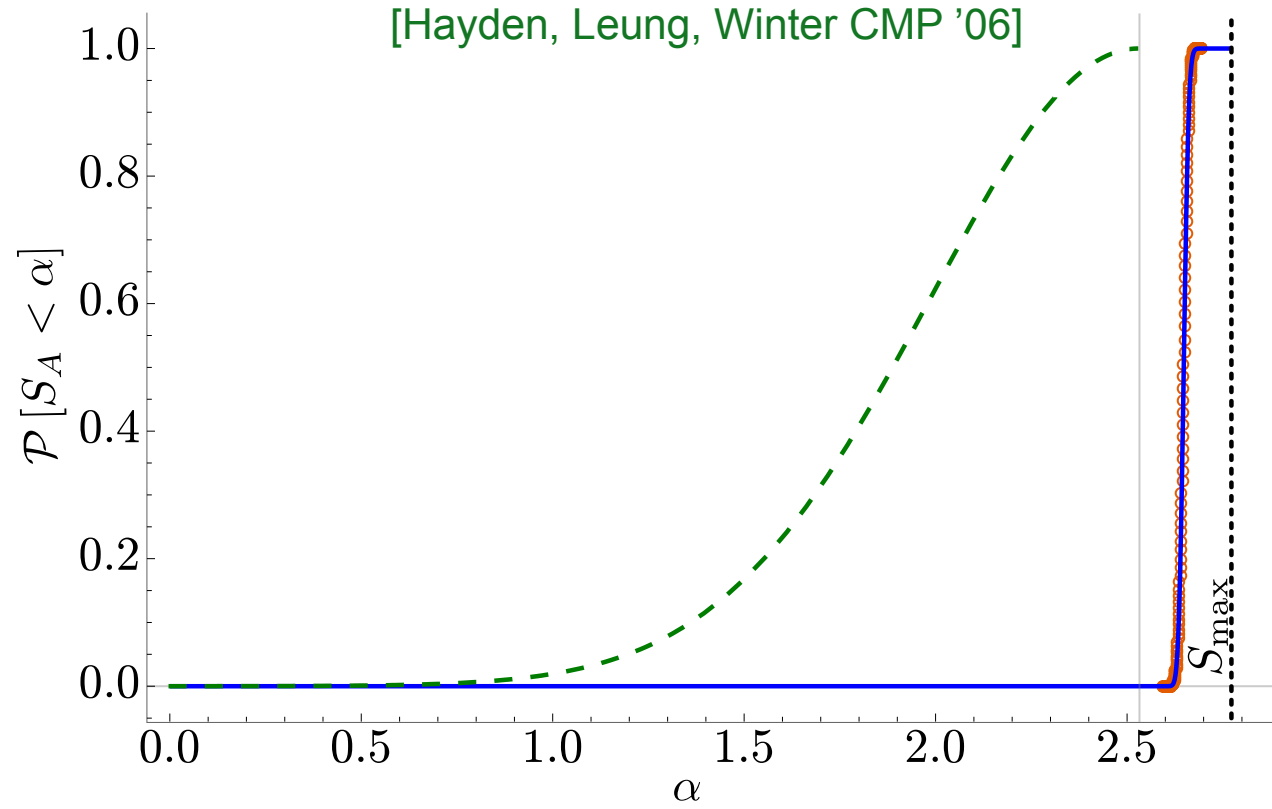
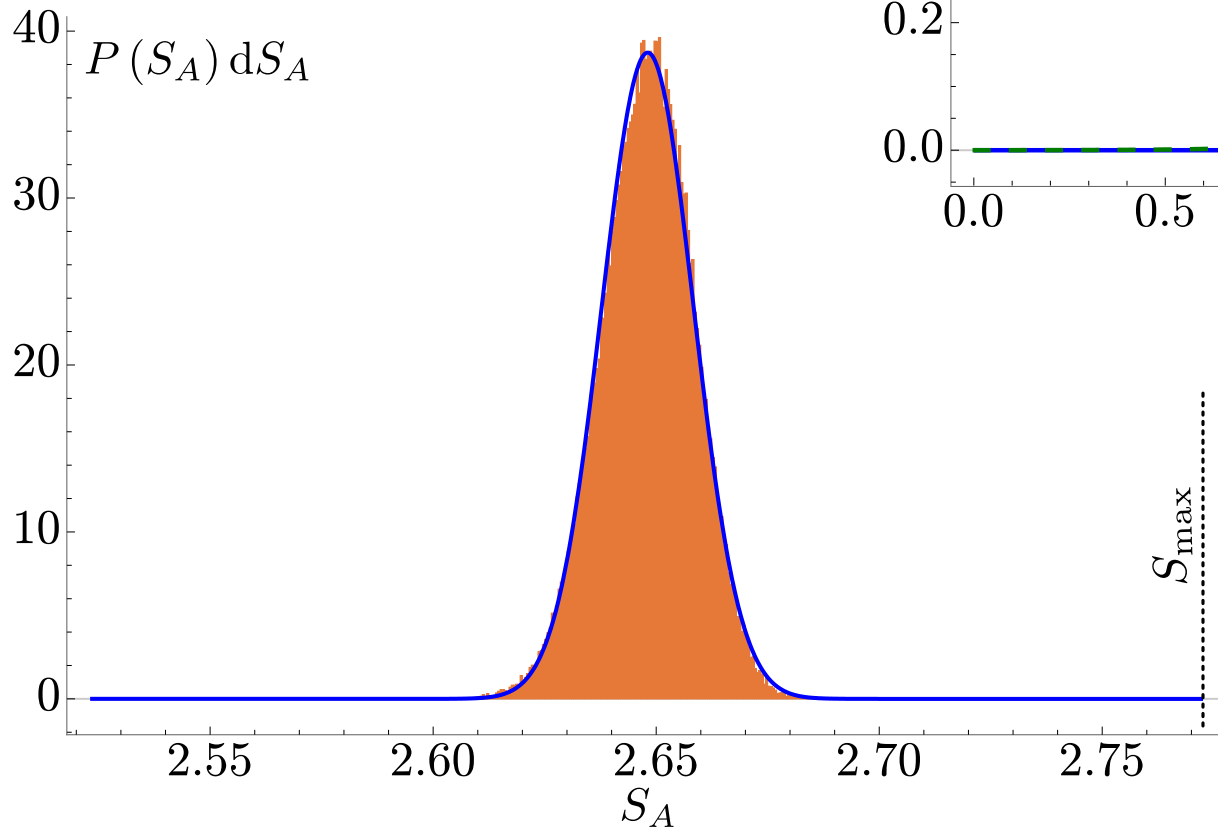
200 bins

$N_A = 4$ $N_B = 6$

Concentration of measure bounds

$$\mathcal{P}[S_A < \alpha] \leq e^{-\frac{(\alpha - \mu_b)^2}{2\sigma_b^2}}$$

$$\mu_b = \log d_A - \frac{d_A}{2d_B} \quad \sigma_b = \frac{2\pi \log d_A}{\sqrt{d_A d_B - 1}}$$



Hausdorff's moment problem
Any moment is computable (skewness)
Normal distribution is a good approximation

[Bianchi, P.D. PRD 2019]

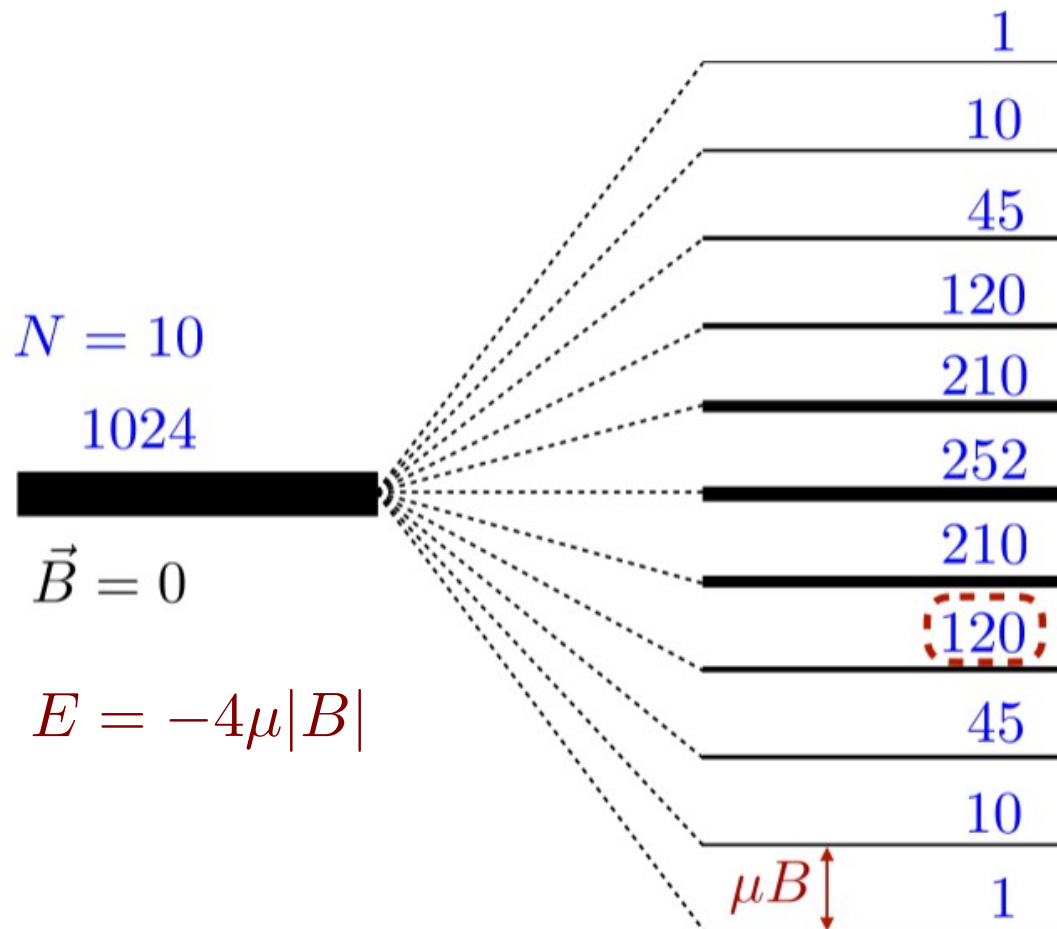
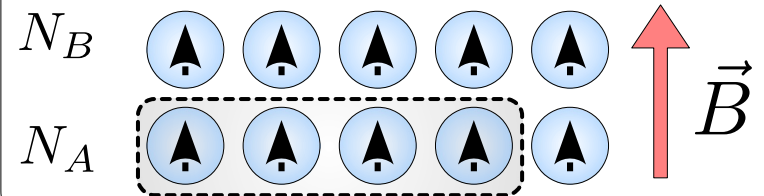
Typical entropy of a subsystem with a center (example)

Hamiltonian $H = \sum_i \mu \vec{\sigma}_i \cdot \vec{B}$

eigenspace with
given energy
(direct sum)

$$\mathcal{H}^E = \bigoplus_{\epsilon} \left(\mathcal{H}_A^{(\epsilon)} \otimes \mathcal{H}_B^{(\epsilon)} \right)$$

N spins + magn. field



Paramagnetic ionic salt

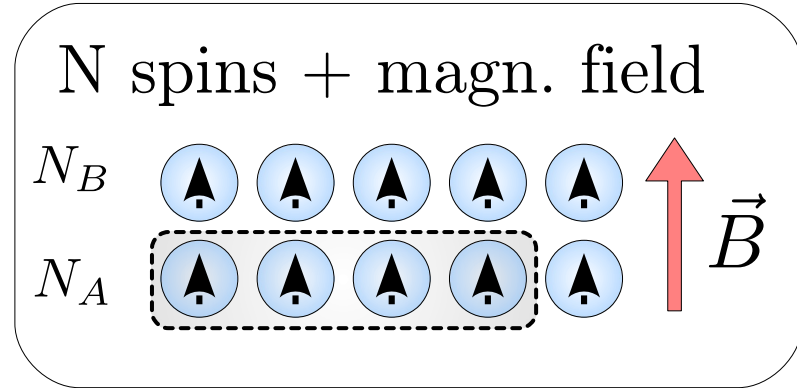


$\mu \simeq 0.9 \times 10^{-23} \text{ J/T}$

$\simeq 0.6 \times 10^{-4} \text{ eV/T}$

Typical entropy of a subsystem with a center (Example)

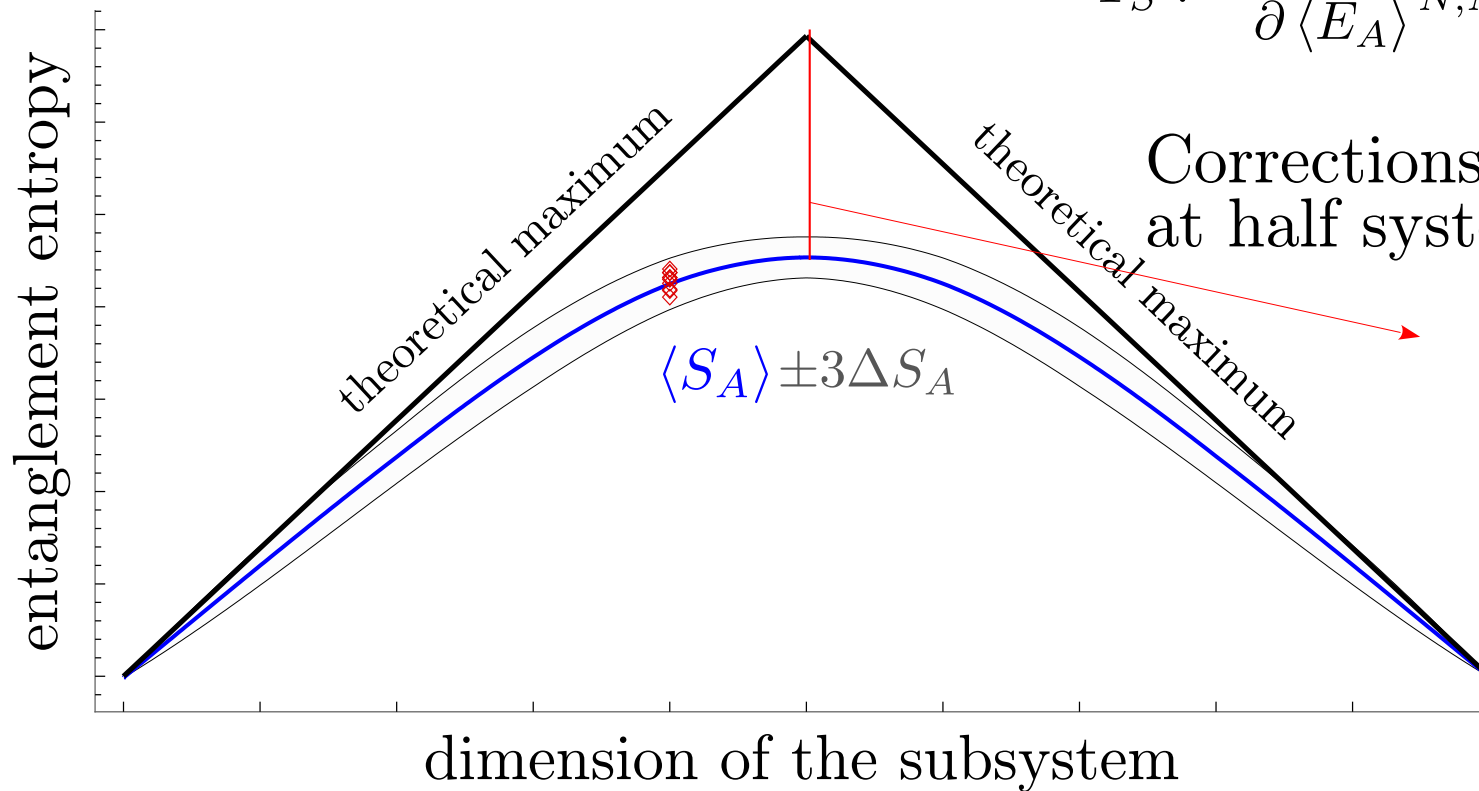
$$\langle S_A \rangle = -\frac{N_A}{2} \left(1 - \frac{E}{\mu B N} \right) \log \left(\frac{1 - E/(\mu B N)}{2} \right) - \frac{N_A}{2} \left(1 + \frac{E}{\mu B N} \right) \log \left(\frac{1 + E/(\mu B N)}{2} \right)$$



The average entropy is typical
[Bianchi, P.D. PRD 2019]

Temperature from entanglement

$$T_S := \frac{\partial \langle S_A \rangle}{\partial \langle E_A \rangle}_{N, N_A} = \frac{\mu B}{\operatorname{arctanh} E/(\mu B N)}$$



Corrections to entanglement
at half system (heat capacity)

$$\Delta S_{half} = \pi \sqrt{N} \sqrt{C}$$

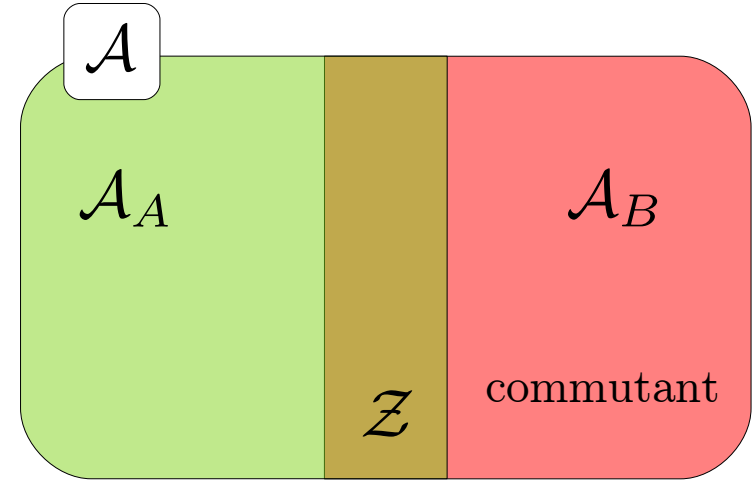
[Murthy, Srednicki PRE 2019,
Bianchi, P.D. wip]

Typical entropy in presence of a center [Bianchi, P.D. PRD 2019]

Hilbert space structure:

$$\mathcal{H} = \bigoplus_{\zeta \in \mathcal{Z}} \left(\mathcal{H}_A^{(\zeta)} \otimes \mathcal{H}_B^{(\zeta)} \right)$$

$$|\psi\rangle = \sum_{\zeta} \sqrt{p_{\zeta}} |\phi_A^{(\zeta)}\rangle |\phi_B^{(\zeta)}\rangle$$



Entanglement entropy

$$S_A(\psi) = \sum_{\zeta} p_{\zeta} S_A(|\phi_A^{(\zeta)}\rangle |\phi_B^{(\zeta)}\rangle) - \sum_{\zeta} p_{\zeta} \log p_{\zeta}$$

The average entropy is typical

$$\langle S_A(\psi) \rangle = \sum_{\zeta} \frac{d_{A\zeta} d_{B\zeta}}{d} \langle S_{A\zeta} \rangle + \Psi(d+1) - \Psi(d_{A\zeta} d_{B\zeta} + 1)$$

Exact formula. Variance and other momenta also computed.

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Information-theoretic bounds on correlations

State:

$$|\psi\rangle \in \mathcal{H}_\Gamma$$

Subsystem: A

$$O_A \in \mathcal{A}_A$$

Correlations

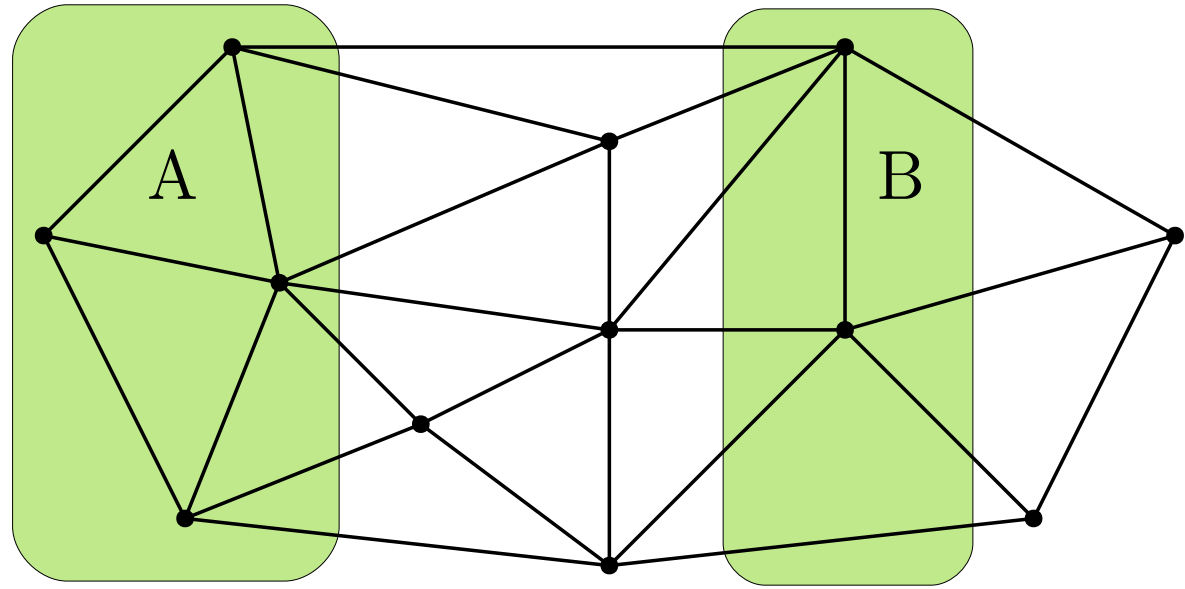
$$\mathcal{C}(O_A, O_B) = \langle \psi | O_A O_B | \psi \rangle - \langle \psi | O_A | \psi \rangle \langle \psi | O_B | \psi \rangle$$

Bounded by mutual information

$$\frac{1}{2} \left(\frac{\mathcal{C}(O_A, O_B)}{\|O_A\| \|O_B\|} \right)^2 \leq S_A(\psi) + S_B(\psi) - S_{AB}(\psi)$$

[Wolf, Verstraete, Hastings, Cirac PRL 2008]

random states (large dimension) follow a volume law
= no correlations



Bell-network states

Gluing quantum polyhedra with entanglement

[Bianchi, Baytas, Yokomizo, PRD 2018]

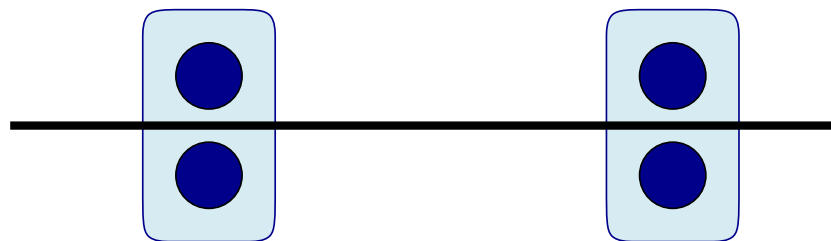
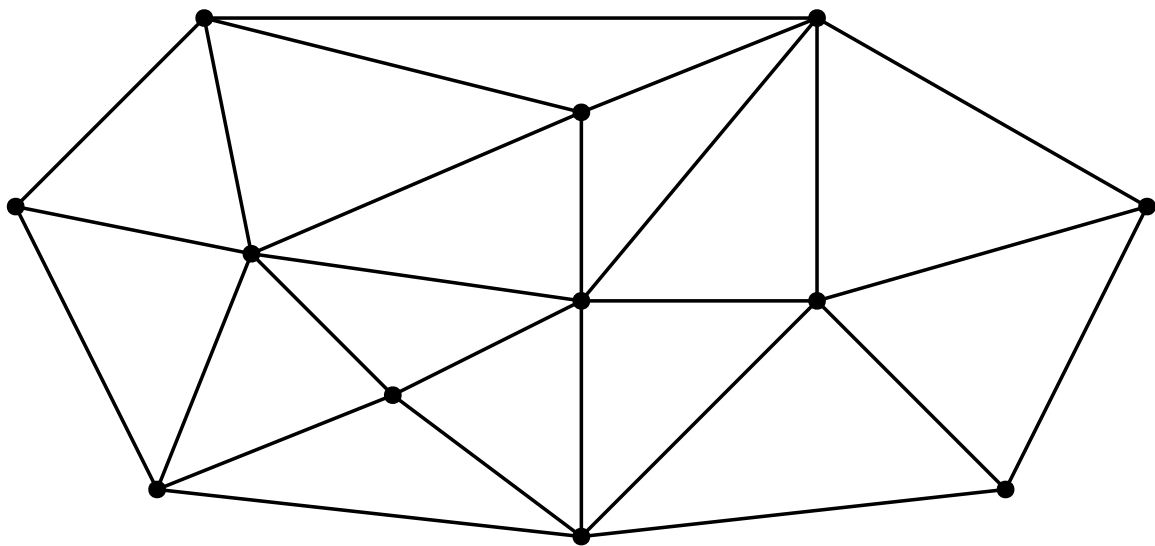
fluctuations of nearby quantum polyhedra are in general uncorrelated (twisted geometry)

Use squeezed vacua techniques to correlate shapes link by link

[Bianchi, Hackl, Guglielmon, Yokomizo, PRD 2016]

Enlarge Hilbert space, squeeze the oscillators and project into \mathcal{H}_Γ

$$\mathcal{H}_\Gamma = \bigoplus_{j_\ell} \bigotimes_n \mathcal{H}_n$$



Schwinger rep of SU(2)

$$|\Gamma, \lambda_\ell, \mathcal{B}\rangle = \sum_{j_\ell} \prod_{\ell} (1 - |\lambda_\ell|^2) \lambda_\ell^{2j_\ell} \sqrt{2j_\ell + 1} |\Gamma, j_\ell, \mathcal{B}\rangle \quad |\Gamma, j_\ell, \mathcal{B}\rangle = \sum_{i_n} \mathcal{S}_\Gamma(j_\ell, i_n) \bigotimes_n |i_n\rangle$$

$\lambda_\ell \in \mathbb{C}$ link squeezing parameter (average area, extrinsic curvature)

symbol of the graph

Bell-network states (analytic and numerics)

$$\mathcal{H}_\Gamma = \bigoplus_{j_\ell} \bigotimes_n \mathcal{H}_n$$

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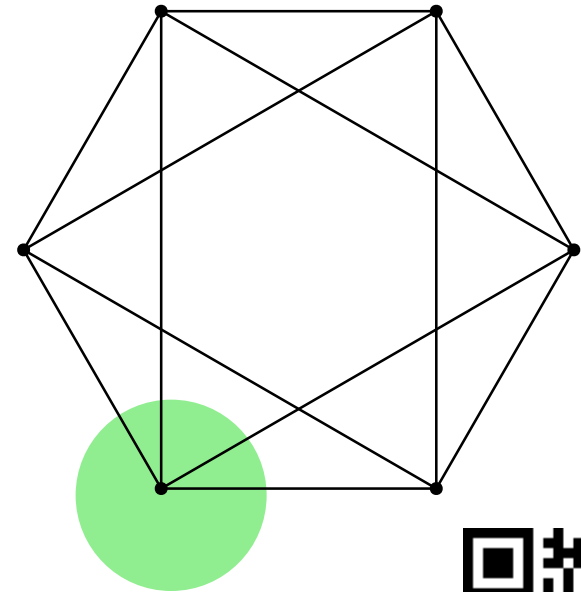
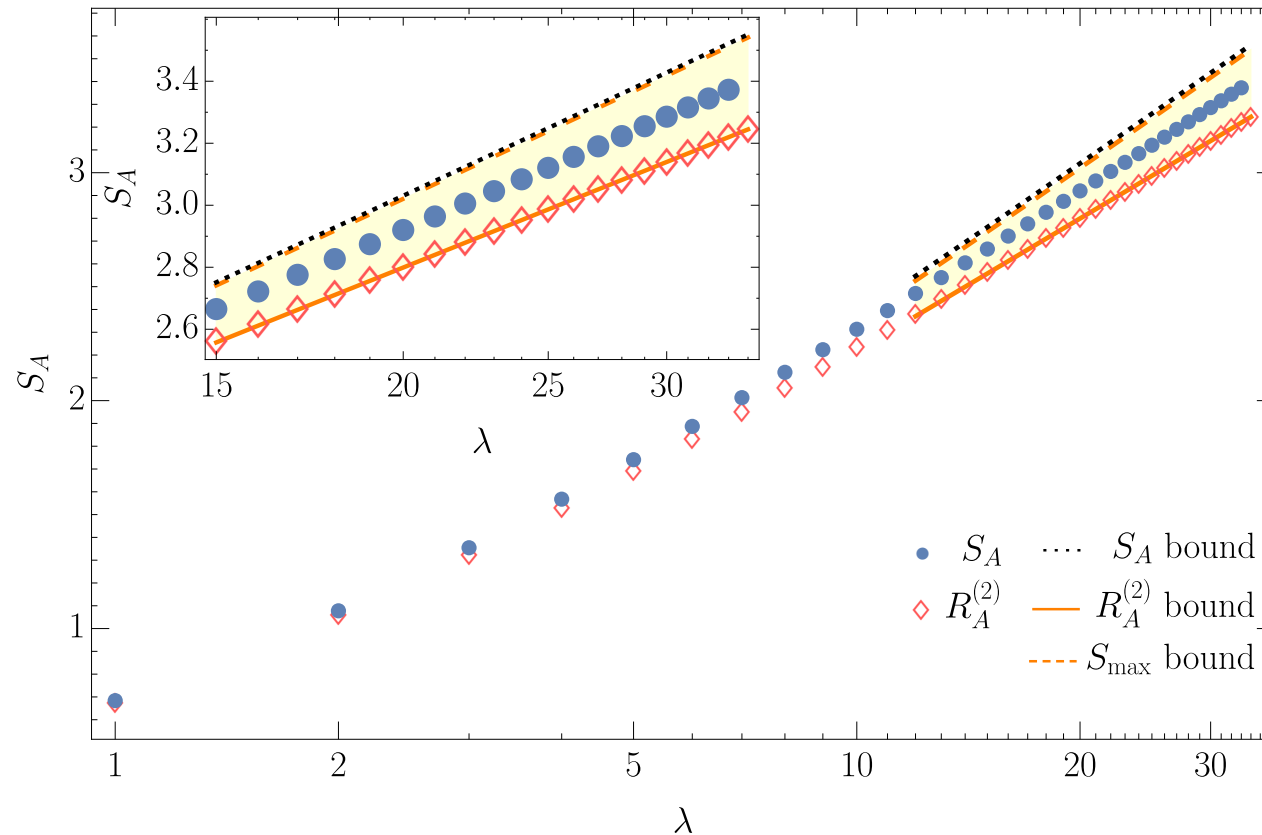
[Bianchi, P.D, Vilensky PRD 2019]

Asymptotic bound on the entanglement entropy

$$j_\ell \rightarrow \lambda j_\ell$$

$$\left(|\partial A| - \frac{C_{\Gamma,A}}{2} \right) \log \lambda \leq S_A \leq \left(|\partial A| - \frac{3}{2} \right) \log \lambda$$

Numeric computation is available for now only on small graphs (but large spins)



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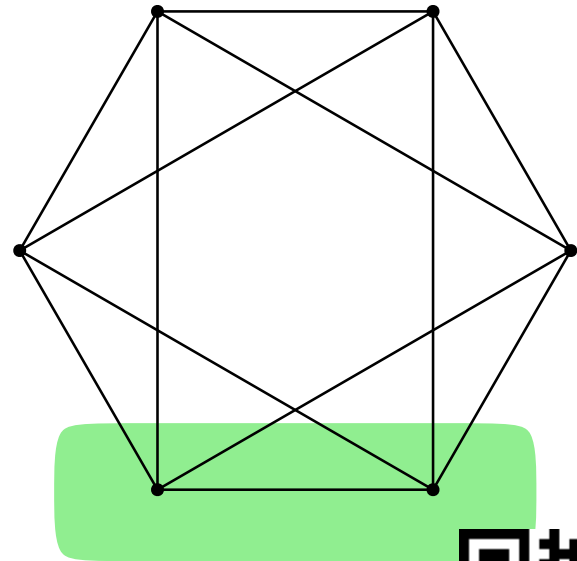
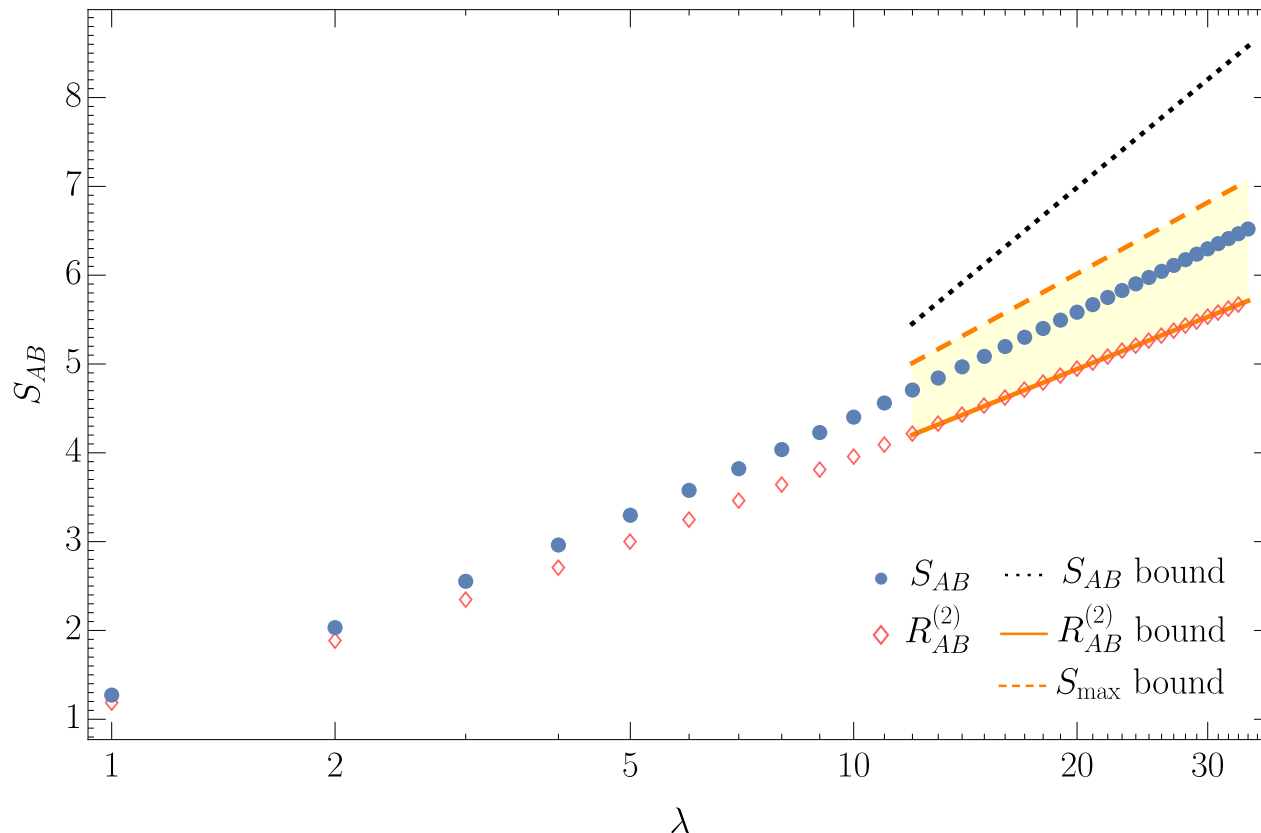
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From numerical fit we can infer the bound on correlations

$$\frac{1}{2} \left(\frac{\mathcal{C}(O_A \cdot O_B)}{\|O_A\| \|O_B\|} \right)^2 \leq 0.06 \log \lambda$$

while for random states we obtain

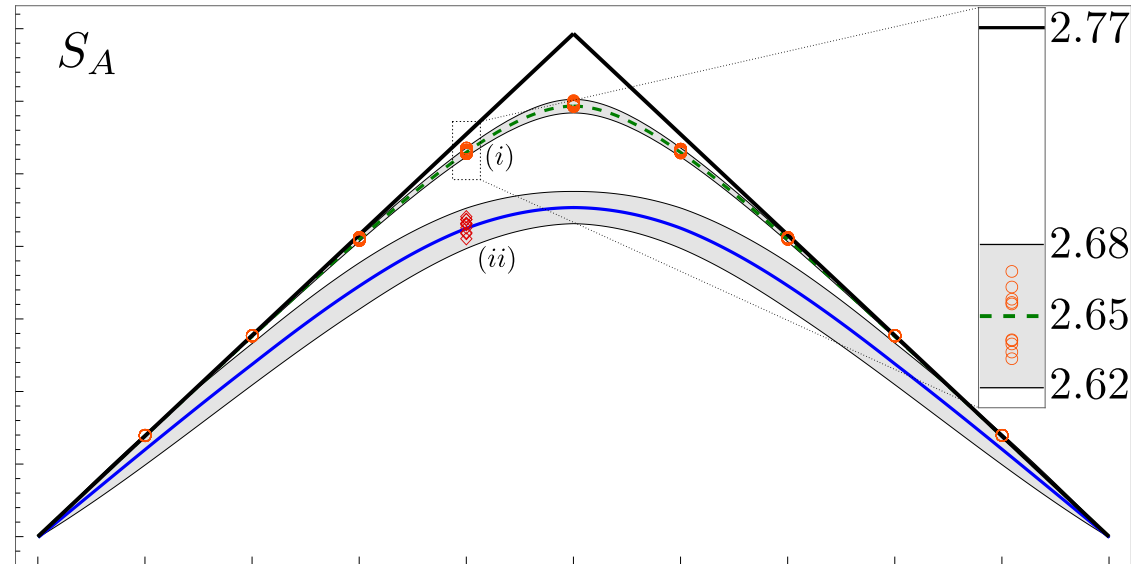
$$\frac{1}{2} \left(\frac{\mathcal{C}(O_A \cdot O_B)}{\|O_A\| \|O_B\|} \right)^2 \leq O(\lambda^{-1})$$



Conclusions

Page curve and its variance:

- Random states have typical entropy
- Unlikely to have maximum entropy
- Concentration of measure
- Vanishing correlations
- Even in presence of a center
- Temperature from entanglement
- Half volume correction (heat capacity)



Entanglement entropy in a Bell-network state:

- Analytic asymptotics
- Area-law from intertwiner entanglement
- Numerical code
- Non-vanishing intertwiner correlations

